# MATHEMATICS OF COILING IN CYLINDRICAL ELECTROCHEMICAL CELLS: THE THEORY OF A SPIRAL BOUNDED BY TWO CIRCLES AND ITS APPLICATION TO THE SPIRAL-WOUND NICKEL-CADMIUM CELL 

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Summary
The mathematics relating to a spiral bounded by two circles has been developed and then applied to the spiral-wound nickel-cadmium cell. Given the input parameters, internal can diameter, component thicknesses and mandrel radius the maximum length of each component may be calculated for 16 different variations.

## 1. Introduction

A coil pack in an electrochemical cell is usually wound on a cylindrical mandrel and fits into a can with circular cross-section. A section normal to the axis of either the mandrel or the can therefore shows each component following the locus of a spiral. In order to maximize cell capacity the gap between the spiral-wound pack and the can should, as far as possible, be filled with active material by overlapping the electrodes.

The relevant mathematics of a spiral bounded by two circles is presented in this paper. This is then applied to the coil pack of the nickelcadmium cell so that electrode overlap and maximum component lengths may be determined.

## 2. General mathematics

The equation for a spiral (Fig. 1) as given by Marincic [1] is:

$$
\begin{equation*}
\rho_{n}=r+n t \tag{1}
\end{equation*}
$$

$\rho_{n}$ is the distance of the spiral from the centre. The spiral originates on the circumference of a circle radius $r$ and moves out a distance $t$ for each turn; $n$ is the number of turns. The length of the spiral ( $L$ ) is given by [1]:


Fig. 1. Geometry of the spiral.

$$
\begin{equation*}
L=\pi\left(n^{2} t+2 r n\right) \tag{2}
\end{equation*}
$$

Marincic [1] states that the minimal diameter $\left(D_{c}\right)$ of the circle containing the spiral is given by:

$$
\begin{equation*}
D_{c}=\rho_{n}+\rho_{(n-0.5)} \tag{3}
\end{equation*}
$$

However, it is shown in Fig. 2 that this is not strictly correct. The diameter given by eqn. (3) is shown as AB . The true minimum diameter $\left(D_{t}\right)$ is shown as AC. Let the number of turns to $C$ and $A$ be $n^{\prime}$ and $n$ respectively. Angles $\angle O A C$ and $\angle A C O$ are shown in Fig. 2 as $\epsilon$ and $\beta$ respectively. It is shown in the Appendix that:

$$
\begin{equation*}
\tan \beta=t / 2 \pi\left(r+n^{\prime} t\right) \tag{4}
\end{equation*}
$$

Using eqn. (1):

$$
\begin{equation*}
\mathrm{OP}=\left(r+n^{\prime} t\right) \sin \beta=(r+n t) \sin \epsilon \tag{5}
\end{equation*}
$$

Also

$$
\begin{equation*}
n-n^{\prime}=\{\pi-(\epsilon+\beta)\} / 2 \pi \tag{6}
\end{equation*}
$$

Rearranging eqn. (4):


Fig. 2. Envelopment of a spiral by a circle. E is a point on the spiral. Straight lines OE and DE intersect the enveloping circle at points $J$ and $F$ respectively.

$$
\begin{equation*}
n^{\prime}=\frac{1}{2 \pi \tan \beta}-\frac{r}{t} \tag{7}
\end{equation*}
$$

From eqn. (5):

$$
\begin{equation*}
\epsilon=\sin ^{-1} \frac{\left(r+n^{\prime} t\right) \sin \beta}{(r+n t)} \tag{8}
\end{equation*}
$$

From eqn. (6):

$$
\begin{equation*}
\frac{n-1 / 2+(\epsilon+\beta) / 2 \pi}{n^{\prime}}-1=0 \tag{9}
\end{equation*}
$$

Equations (7) - (9) were solved by an iterative method using a Texas SR-52 programmable calculator. For particular values of $n, r$ and $t$ a value was assumed for $\beta$ and $n^{\prime}$ calculated using eqn. (7). Then $\epsilon$ was evaluated using eqn. (8) and hence the l.h.s. of eqn. (9) was computed. $\beta$ was adjusted until the l.h.s. of eqn. (9) was very small. The true value for the minimum diameter of the enveloping circle $\left(D_{t}\right)$ was then calculated according to eqn. (10) and also the fractional error $\left(D_{t} / D_{c}-1\right)$ made in using eqn. (3):

$$
\begin{equation*}
D_{t}=\mathrm{AP}+\mathrm{PC}=\left(r+n^{\prime} t\right) \cos \beta+(r+n t) \cos \epsilon \tag{10}
\end{equation*}
$$

Values obtained when $\{n-1 / 2+(\epsilon+\beta) / 2 \pi\} n^{\prime}-1<10^{-7}$ are presented in Table 1. The maximum error resulting from the use of eqn. (3)

TABLE 1
Computations of the error resulting from the use of eqn. (3)

| $n$ |  | $t / r$ |  |  |  |  |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- |
|  |  | 0.1 | 0.3 | 1 | 3 | 10 |
| 1 | $n^{\prime}$ | 0.505 | 0.512 | 0.529 | 0.547 | 0.560 |
|  | $D_{t} / r$ | 2.150 | 2.452 | 3.515 | 6.572 | 17.31 |
|  | $\epsilon^{\circ}$ | 0.83 | 2.10 | 4.54 | 6.75 | 8.09 |
|  | $\beta^{\circ}$ | 0.87 | 2.37 | 5.94 | 10.25 | 13.56 |
|  | $\left(D_{t} / D_{c}-1\right)$ | $1.1 \times 10^{-4}$ | $7.6 \times 10^{-4}$ | $4.2 \times 10^{-3}$ | $1.1 \times 10^{-2}$ | $1.8 \times 10^{-2}$ |
| 3 | $n^{\prime}$ |  |  |  |  |  |
|  | $D_{t} / r$ | 2.504 | 2.508 | 2.514 | 2.516 | 2.518 |
|  | $\epsilon^{\circ}$ | 0.70 | 3.651 | 7.507 | 18.525 | 57.09 |
|  | $\beta^{\circ}$ | 0.73 | 1.44 | 2.28 | 2.73 | 2.94 |
|  | $\left(D_{t} / D_{c}-1\right)$ | $7.8 \times 10^{-5}$ | 1.56 | $3.4 \times 10^{-4}$ | 2.59 | $9.0 \times 10^{-4}$ |
| 10 | $n^{\prime}$ | 9.503 | 9.504 | 9.20 | $1.3 \times 10^{-3}$ | 3.48 |
|  | $D_{t} / r$ | 3.950 | 7.851 | 21.505 | $9.50^{-3}$ |  |
|  | $\epsilon^{\circ}$ | 0.46 | 0.68 | 0.83 | 60.51 | 197.0 |
|  | $\beta^{\circ}$ | 0.47 | 0.71 | 0.87 | 0.88 | 0.90 |
|  | $\left(D_{t} / D_{c}-1\right)$ | $3.2 \times 10^{-5}$ | $7.4 \times 10^{-5}$ | $1.1 \times 10^{-4}$ | $1.2 \times 10^{-4}$ | $1.3 \times 10^{-4}$ |

is $1.8 \%$ for $n=1$ and $t / r=10$ which is an extreme case. In the nickelcadmium cell typical values are $n=7$ and $t / r=0.5$ for which the calculated error in using eqn. (3) is $<0.02 \%$. Equation (3) is therefore an extremely close approximation and will be used subsequently for the calculation of $n$.

It can be seen in Table 1 that $\epsilon$ and $\beta$ are usually small and hence the approximation $\beta=\sin \beta=\tan \beta, \epsilon=\sin \epsilon$ for angles in radians may be made to allow eqns. (4) - (6) to be solved. Equations (4) and (5) then become:

$$
\begin{align*}
& \beta=t / 2 \pi\left(r+n^{\prime} t\right)  \tag{11}\\
& \left(r+n^{\prime} t\right) \beta=(r+n t) \epsilon \tag{12}
\end{align*}
$$

Equating ( $\left.r+n^{\prime} t\right) \beta$ in eqns. (11) and (12):

$$
\begin{equation*}
\epsilon=t / 2 \pi(r+n t) \tag{13}
\end{equation*}
$$

Substituting for $\epsilon$ and $\beta$ from eqns. (11) and (13) into eqn. (6) and writing $2 \pi(n-1 / 2)+t / 2 \pi(r+n t)=K$ leads to:

$$
\begin{equation*}
4 \pi^{2} t\left(n^{\prime}\right)^{2}+\left(4 \pi^{2} r-2 \pi K t\right) n^{\prime}-(t+2 \pi K r)=0 \tag{14}
\end{equation*}
$$

This is a quadratic equation in $n^{\prime}$. Hence given $n, r$ and $t, n^{\prime}$ may be calculated from eqn. (14) and $\beta$ and $\epsilon$ may be determined from eqns. (11) and (13).

## 3. Distances between the spiral and the outer circle

### 3.1. Theory

The required distances EF and EJ for the determination of possible electrode overlap are shown in Fig. 2. The centre of the enveloping circle is D. The essential parts of Fig. 2 required for a geometric analysis are shown in Fig. 3. It should be noted that $\mathrm{OC}=\rho_{n^{\prime}}, \mathrm{OA}=\rho_{n}$ and $\mathrm{JD}=\mathrm{CD}=D_{t} / 2$.

The angles $\lambda, \beta, \epsilon, \alpha, \theta$ and $\gamma$ are shown in Fig. 3 and all other angles may be expressed in terms of these six. It can be seen using eqn. (1) that:

$$
\begin{equation*}
\mathrm{OE}=r+\{(n-1)+(\lambda-\alpha) / 2 \pi\} t \tag{15}
\end{equation*}
$$

Furthermore

$$
\begin{align*}
& \mathrm{OG}=\mathrm{OD} \cos (\lambda-\theta)  \tag{16}\\
& \mathrm{GE}=\mathrm{OD} \sin (\lambda-\theta) \cot (-\lambda+\alpha+\epsilon+\gamma) \tag{17}
\end{align*}
$$

But OE $=\mathrm{OG}+\mathrm{GE}$. Therefore from eqns. (15) - (17):

$$
\begin{gather*}
r+\{(n-1)+(\lambda-\alpha) / 2 \pi\} t=\mathrm{OD}\{\cos (\lambda-\theta)+\sin (\lambda-\theta) \cot (-\lambda+ \\
\alpha+\epsilon+\gamma)\} \tag{18}
\end{gather*}
$$



Fig. 3. The essential parts of Fig. 2 for a mathematical analysis.

$$
\begin{align*}
& \mathrm{OH}=\rho_{n^{\prime}} \sin \beta=\mathrm{OD} \sin (\alpha+\epsilon-\theta)  \tag{19}\\
& \mathrm{OD}=\left(\rho_{n^{\prime}} \sin \beta\right) / \sin (\alpha+\epsilon-\theta) \tag{20}
\end{align*}
$$

But $\mathrm{CH}+\mathrm{HD}=D_{t} / 2$ and therefore

$$
\begin{align*}
& \rho_{n^{\prime}} \cos \beta+\left(\rho_{n^{\prime}} \sin \beta\right) \cot (\alpha+\epsilon-\theta)=D_{t} / 2  \tag{21}\\
& \theta=\alpha+\epsilon-\tan ^{-1} \frac{\rho_{n^{\prime}} \sin \beta}{D_{t} / 2-\rho_{n^{\prime}} \cos \beta} \tag{22}
\end{align*}
$$

Also

$$
\begin{equation*}
\mathrm{ED}=\mathrm{OD} \sin (\lambda-\theta) / \sin (-\lambda+\alpha+\epsilon+\gamma) \tag{23}
\end{equation*}
$$

Next an expression for OJ will be derived:

$$
\begin{equation*}
\angle \mathrm{OJD}=\sin ^{-1} \frac{\mathrm{OD} \sin (\lambda-\theta)}{\mathrm{JD}} \tag{24}
\end{equation*}
$$

But OJ = OG + GJ and therefore:

$$
\begin{equation*}
\mathrm{OJ}=\mathrm{OD} \cos (\lambda-\theta)+\frac{D_{t}}{2} \cos \left\{\sin ^{-1} \frac{\mathrm{OD} \sin (\lambda-\theta)}{D_{t} / 2}\right\} \tag{25}
\end{equation*}
$$

The above equations then allow the determination of the required quantities EF and EJ for given values of $D_{c}, r, t$ and $\gamma$ as follows: (i) calculate $n$ using eqns. (1) and (3); (ii) calculate $n^{\prime}$ by solving eqn. (14); (iii) calculate $\beta$ and $\epsilon$ using eqns. (11) and (13); (iv) calculate $D_{t}$ using eqn. (10); (v) calculate $\alpha$ from the decimal part of $n$; (vi) calculate $\theta$ using eqns. (22) and (1); (vii) calculate OD using eqns. (20) and (1); (viii) use an iterative method for solving eqn. (18) i.e. insert a $\lambda$ value and calculate the l.h.s. and r.h.s. of eqn. (18). Adjust $\lambda$ until (l.h.s. - r.h.s.) is very small; (ix) calculate EF = $D_{t} / 2-\mathrm{ED}$ using eqn. (23); (x) calculate $\mathrm{EJ}=\mathrm{OJ}-\mathrm{OE}$ using eqns. (25) and (15). Steps (i) - (x) can be performed on a programmable calculator. In all the calculations which were done the more positive root of eqn. (14) was the correct one and hence the other root may be ignored.

### 3.2. Results

Calculated values of $\lambda-\alpha, \mathrm{EF} / t$ and $\mathrm{EJ} / t$ are shown in Table 2 for values of $D_{t} / t$ from 2.34-43.8 and $\gamma$ from 0-180 .

A plot of EJ/t vs. $\lambda-\alpha$ is shown in Fig. 4. Little dependence on $D_{t} / t$ is indicated. A linear regression analysis was performed on the data $0^{\circ} \leqslant$ $\lambda-\alpha<121^{\circ}$ resulting in the equation (for angles in degrees):

$$
\begin{equation*}
\mathrm{EJ} / t=1.0027-0.007078(\lambda-\alpha) \tag{26}
\end{equation*}
$$

The coefficient of determination was 0.998 indicating that the points correlate well with the straight line in the range covered. The line given by eqn. (26) is shown broken in Fig. 4.

TABLE 2
The distances EF and EJ between the spiral and the outer circle (angles are in degrees)

| $D_{t} / t$ | $\gamma$ | $\lambda-\alpha$ | EF/t | EJ/t | $D_{t} / t$ | $\gamma$ | $\lambda-\alpha$ | EF/t | $\mathrm{EJ} / \mathrm{t}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 2.3437 | 0 | - | 1.08398 | - | 7.8879 | 0 | - | 1.00285 | - |
|  | 30 | - | 0.97471 | - |  | - | 0 | - | 1 |
|  | - | 0 | - | 1 |  | 30 | 27.62 | 0.81839 | 0.82044 |
|  | 60 | 32.78 | 0.73210 | 0.78027 |  | 60 | 57.20 | 0.59106 | 0.59296 |
|  | 90 | 74.63 | 0.43990 | 0.46023 |  | 90 | 88.11 | 0.35803 | 0.35885 |
|  | 120 | 117.09 | 0.19358 | 0.19576 |  | 120 | 120.04 | 0.16257 | 0.16268 |
|  | 150 | 157.96 | 0.04434 | 0.04435 |  | 150 | 152.43 | 0.03919 | 0.03919 |
|  | 180 | 196.01 | 0 | 0 |  | 180 | 184.64 | 0 | 0 |
| 2.6726 | 0 | - | 1.04640 | - | 20.057 | 0 | - | 1.00040 | - |
|  | - | 0 | - | 1 |  | - | 0 | - | 1 |
|  | 30 | 11.12 | 0.90607 | 0.93321 |  | 30 | 29.29 | 0.80834 | 0.80863 |
|  | 60 | 41.91 | 0.68133 | 0.71081 |  | 60 | 59.10 | 0.57759 | 0.57785 |
|  | 90 | 78.81 | 0.41684 | 0.43966 |  | 90 | 89.40 | 0.34700 | 0.34712 |
|  | 120 | 118.04 | 0.18656 | 0.18803 |  | 120 | 120.08 | 0.15706 | 0.15708 |
|  | 150 | 156.99 | 0.04340 | 0.04341 |  | 150 | 150.97 | 0.03800 | 0.03799 |
|  | 180 | 193.96 | 0 | 0 |  | 180 | 181.82 | 0 | 0 |
| 4.0128 | 0 | - | 1.01387 | - | 43.751 | 0 | - | 1.00008 | - |
|  | - | 0 | - | 1 |  | - | 0 | - | 1 |
|  | 30 | 22.78 | 0.84505 | 0.85456 |  | 30 | 29.70 | 0.80552 | 0.80568 |
|  | 60 | 52.24 | 0.62216 | 0.63136 |  | 60 | 59.62 | 0.57361 | 0.57367 |
|  | 90 | 84.83 | 0.38098 | 0.38500 |  | 90 | 89.75 | 0.34358 | 0.34361 |
|  | 120 | 119.46 | 0.17298 | 0.17349 |  | 120 | 120.05 | 0.15528 | 0.15529 |
|  | 150 | 154.69 | 0.04122 | 0.04122 |  | 150 | 150.45 | 0.03759 | 0.03759 |
|  | 180 | 189.19 | 0 | 0 |  | 180 | 180.83 | 0 | 0 |

Plots of EF/t vs. $D_{t} / t$ are shown in Fig. 5 for values of $\gamma$ from $0-150^{\circ}$. A considerable dependence on $D_{t} / t$ is shown in this case. In order to determine EF/t at low $D_{t} / t$ values it was found convenient to plot EF/t vs. ( $D_{t} / t-$ $2^{-1}$ as shown in Fig. 6. This resulted in smooth curves from which intermediate points could be ascertained more precisely than from plots of $D_{t} / t$ vs. EF/t (Fig. 5). Using the family of curves of the type in Fig. 6 for $\gamma=0-150^{\circ}$ the curves shown in Fig. 7 were plotted of EF/t vs. $\gamma$ for $D_{t} / t$ values of $2.35,3.00$ and 40.0. Each curve approximates to a straight line in the range $30 \leqslant \gamma \leqslant 120$. Linear regression analyses were performed on the data in Table 2 over this range. The following equation may be written:

$$
\begin{equation*}
\mathrm{EF} / t=A \gamma+B(30 \leqslant \gamma \leqslant 120) \tag{27}
\end{equation*}
$$

$A$ and $B$ values are listed in Table 3. These values are plotted vs. $\left(D_{t} / t-2\right)^{-1}$ in Fig. 8. As can be seen the points for $A$ vs. $\left(D_{t} / t-2\right)^{-1}$ lie close to a straight line. A linear regression analysis revealed the relationship:

$$
\begin{equation*}
A=-0.000530\left(D_{t} / t-2\right)^{-1}-0.00726 \tag{28}
\end{equation*}
$$




Fig. 4. EJ/t us. $\lambda-\alpha . D_{t} / t$ values: •, 2.34; $\circ$, 2.67; $\llcorner, 4.01 ; \Delta, 7.89 ; \bullet, 20.1 ; \nabla, 43.8$.
Broken line represents the best straight line through the points $0^{\circ} \leqslant \lambda-\alpha<121^{\circ}$.
Fig. 5. EF/t us. $D_{t} / t$ at the various values of $\gamma$ indicated on the curves.


Fig. 6. EF/t vs. $\left(D_{t} / t-2\right)^{-1}$ for $\gamma=60^{\circ}$.
Fig. 7. EF/t $v s$. $\gamma$ for the following $D_{t} / t$ values: $\qquad$ 40.0

Hence for $D_{t} / t$ values in the range 2.34-43.8, $A$ may be determined from eqn. (28) and $B$ by graphical interpolation from Fig. 8 for substitution into eqn. (27).

TABLE 3
Coefficients of eqn. (27)

| $D_{t} / t$ | $A\left(\right.$ degree $\left.^{-1}\right)$ | $B$ |
| :--- | :--- | :--- |
| 2.3437 | -0.00879 | 1.244 |
| 2.6726 | -0.00808 | 1.153 |
| 4.0128 | -0.00752 | 1.070 |
| 7.8879 | -0.00733 | 1.033 |
| 20.057 | -0.00728 | 1.019 |
| 43.751 | -0.00727 | 1.015 |

Fig. 8. $A$ and $B$ vs. $\left(D_{t} / t-2\right)^{-1}$.

## 4. The spiral-wound construction

The mandrel on which the pack is wound is usually split to allow the use of a single piece of separator (Fig. 9). The length of separator in the split of the mandrel has not been taken into account in the subsequent calculations. Although a single piece of separator is generally used, it is treated here as two pieces, 1 and 2 , with compressed thicknesses $\delta_{1}$ and $\delta_{2}$ respectively making the analysis more general.


Fig. 9. Winding the components onto a split mandrel.

## 5. Cases with no electrode overlap at the inside or outside of the coil pack

There are four basic cases corresponding to the positive or negative electrode starting at the centre and the positive or negative ending at the outside

positive electrode negative electrode SEPARATOR 1 SEPARATOR 2

Fig. 10. (a) Case 1: positive starts at centre and negative ends outside. (b) Case 2: positive starts at centre and ends outside. (c) Case 3: negative starts at centre and positive ends outside. (d) Case 4: negative starts at centre and ends outside.
of the pack. The can is the negative terminal and hence if the positive finishes at the outside it must be insulated from the can by a layer of separator. The four cases are shown in Fig. 10.

Let the mandrel have a radius $r_{c}$ and the positive, separator 1 , negative and separator 2 have thicknesses $\delta_{p}, \delta_{1}, \delta_{N}$ and $\delta_{2}$ and lengths $L_{p}, L_{1}, L_{N}$ and $L_{2}$ respectively. In this and subsequent calculations it will be assumed that the length of a flat piece of electrode or separator is equal to the length of a plane through the centre of the component when it is bent around the mandrel, i.e. it compresses at the inside and stretches at the outside. Using eqn. (2) for the length of a spiral the following relationships follow:

$$
\begin{align*}
& L_{p}=\pi\left[(n-a)^{2} t+2\left\{r_{c}+\delta_{p} / 2+d\left(\delta_{N}+\delta_{2}\right)\right\}(n-a)\right]  \tag{29}\\
& L_{1}=\pi\left[(n-a)^{2} t+2\left\{r_{c}+\delta_{p}+\delta_{1} / 2+d\left(\delta_{N}+\delta_{2}\right)\right\}(n-a)\right] \tag{30}
\end{align*}
$$

TABLE 4
Values of the coefficients in eqns. (29) - (32), (34), (51), (53)-
(55), (58) and (63)

| Constant | Case |  |  |  |
| :--- | :--- | :--- | :--- | :--- |
|  | 1 | 2 | 3 | 4 |
| $a$ | 0 | 0 | 0 | 1 |
| $d$ | 0 | 0 | 1 | 1 |
| $h$ | 0 | 1 | 0 | 0 |
| $k$ | 1 | 1 | 0 | 1 |

$$
\begin{align*}
& L_{N}=\pi\left[(n-h)^{2} t+2\left\{r_{c}+(1-d)\left(\delta_{p}+\delta_{1}\right)+\delta_{N} / 2\right\}(n-h)\right]  \tag{31}\\
& L_{2}=\pi\left[(n-k)^{2} t+2\left\{r_{c}+(1-d)\left(\delta_{p}+\delta_{1}\right)+\delta_{N}+\delta_{2} / 2\right\}(n-k)\right] \tag{32}
\end{align*}
$$

Values for the coefficients $a, d, h$ and $k$ are given in Table 4. Also

$$
\begin{equation*}
t=\delta_{p}+\delta_{1}+\delta_{N}+\delta_{2} \tag{33}
\end{equation*}
$$

The number of turns $n$ of the component at the outside of the coil pack can be determined by application of eqns. (1) and (3):

$$
\begin{equation*}
n=\frac{2 D_{c}-4\left\{r_{c}+(1-a)\left(\delta_{p}+\delta_{1}\right)+(1-h) \delta_{N}+(1-k) \delta_{2}\right\}+t}{4 t} \tag{34}
\end{equation*}
$$

## 6. Cases with electrode overlap at the centre of the coil pack

It is apparent in Fig. 10 that there is void space between the mandrel and the coil in which further positive or negative electrode could be incorporated.

### 6.1. Cases 1 and 2

The two possibilities for filling the void space are shown in Fig. 11 depending upon whether the positive or negative electrode protrudes into the void space.

Positive filling void space
The distance between the mandrel and the positive electrode $d_{1}$ can be determined using eqn. (1):

$$
\begin{equation*}
d_{1}=\sigma t / 2 \pi \tag{35}
\end{equation*}
$$

$\sigma$ is the angle shown in Fig. 11.
Let the positive protrude into the void space making an angle $\psi$ with the negative. Then the maximum protrusion $\psi_{m}$ is given by (Fig. 11):

(o)
(b)

Fig. 11. Filling of void at centre of pack (a) by positive and (b) by negative. Key as for Fig. 10.

$$
\begin{equation*}
\psi_{m}=\left(2 \pi-\sigma_{m}\right) \tag{36}
\end{equation*}
$$

where (using eqn. 35):

$$
\begin{equation*}
\delta_{2}+\delta_{p}+\delta_{1}=\sigma_{m} t / 2 \pi \tag{37}
\end{equation*}
$$

Eliminating $\sigma_{m}$ from eqns. (36) and (37):

$$
\begin{equation*}
\psi_{m}=2 \pi\left\{1-\left(\delta_{2}+\delta_{p}+\delta_{1}\right) / t\right\} \tag{38}
\end{equation*}
$$

This is equivalent to a length of positive $L_{p}^{\prime}$ where

$$
\begin{align*}
& L_{p}^{\prime}=\left(r_{c}+\delta_{p} / 2\right) \psi_{m} \\
& \text { i.e. } L_{p}^{\prime}=2 \pi\left(r_{c}+\delta_{p} / 2\right)\left\{1-\left(\delta_{2}+\delta_{p}+\delta_{1}\right) / t\right\} \tag{39}
\end{align*}
$$

Similarly the corresponding length of separator 1 is

$$
\begin{equation*}
\left(L_{1}^{\prime}\right)_{p}=2 \pi\left(r_{c}+\delta_{p}+\delta_{1} / 2\right)\left\{1-\left(\delta_{2}+\delta_{p}+\delta_{1}\right) / t\right\} \tag{40}
\end{equation*}
$$

Separator 2 protrudes into the void space but continues to follow a spiral with $\psi_{m}$ corresponding to the $n$ value $n_{m}$ ( $n_{m}$ has a negative value).

$$
\begin{equation*}
n_{m}=-\psi_{m} / 2 \pi=\left(\delta_{2}+\delta_{p}+\delta_{1}\right) / t-1 \tag{41}
\end{equation*}
$$

This corresponds to a length of separator 2 given by:

$$
\begin{align*}
\left(L_{2}^{\prime}\right)_{p}= & -\pi\left[\left\{\left(\delta_{2}+\delta_{p}+\delta_{1}\right) / t-1\right\}^{2} t\right. \\
& \left.+2\left(r_{c}+\delta_{p}+\delta_{1}+\delta_{N}+\delta_{2} / 2\right)\left\{\left(\delta_{2}+\delta_{p}+\delta_{1}\right) / t-1\right\}\right] \tag{42}
\end{align*}
$$

Negative filling the void space
Let the negative protrude into the void space with a maximum length $L_{N}^{\prime}$. The equation equivalent to eqn. (38) is now:

$$
\begin{equation*}
\psi_{m}=2 \pi\left\{1-\left(\delta_{N}+\delta_{2}\right) / t\right\} \tag{43}
\end{equation*}
$$

This leads to ( $c f$. eqn. 42):

$$
\begin{equation*}
L_{N}^{\prime}=-\pi\left[\left\{\left(\delta_{N}+\delta_{2}\right) / t-1\right\}^{2} t+2\left(r_{c}+\delta_{p}+\delta_{1}+\delta_{N} / 2\right)\left\{\left(\delta_{N}+\delta_{2}\right) / t-1\right\}\right] \tag{44}
\end{equation*}
$$

Similarly the corresponding length of separator 2 is:

$$
\begin{align*}
\left(L_{2}^{\prime}\right)_{N}=-\pi\left[\left\{\left(\delta_{N}+\delta_{2}\right) / t-1\right\}^{2} t+2\left(r_{\mathrm{c}}+\delta_{p}+\right.\right. & \left.\delta_{1}+\delta_{N}+\delta_{2} / 2\right) \\
& \left.\left.\left\{\left(\delta_{N}+\delta_{2}\right) / t-1\right)\right\}\right] \tag{45}
\end{align*}
$$

### 6.2. Cases 3 and 4

The treatment is as in the above section except that the positive and separator 1 are interchanged with the negative and separator 2 respectively.

Positive filling void space
The relevant previous equations are (44) and (45):

$$
\begin{equation*}
\left.L_{p}^{\prime}=-\pi\left[\left\{\left(\delta_{p}+\delta_{1}\right) / t-1\right\}^{2} t+2\left(r_{c}+\delta_{N}+\delta_{2}+\delta_{p} / 2\right)\left\{\left(\delta_{p}+\delta_{1}\right) / t-1\right)\right\}\right] \tag{46}
\end{equation*}
$$

$$
\begin{align*}
\left(L_{1}^{\prime}\right)_{p}= & -\pi\left[\left\{\left(\delta_{p}+\delta_{1}\right) / t-1\right\}^{2} t\right. \\
& \left.\left.+2\left(r_{c}+\delta_{N}+\delta_{2}+\delta_{p}+\delta_{1} / 2\right)\left\{\left(\delta_{p}+\delta_{1}\right) / t-1\right)\right\}\right] \tag{47}
\end{align*}
$$

Negative filling the void space
The relevant previous equations are (39), (40) and (42):

$$
\begin{align*}
L_{N}^{\prime}= & 2 \pi\left(r_{c}+\delta_{N} / 2\right)\left\{1-\left(\delta_{1}+\delta_{N}+\delta_{2}\right) / t\right\}  \tag{48}\\
\left(L_{2}^{\prime}\right)_{N}= & 2 \pi\left(r_{c}+\delta_{N}+\delta_{2} / 2\right)\left\{1-\left(\delta_{1}+\delta_{N}+\delta_{2}\right) / t\right\}  \tag{49}\\
\left(L_{1}^{\prime}\right)_{N}= & -\pi\left[\left\{\left(\delta_{1}+\delta_{N}+\delta_{2}\right) / t-1\right\}^{2} t\right. \\
& \left.+2\left(r_{c}+\delta_{N}+\delta_{2}+\delta_{p}+\delta_{1} / 2\right)\left\{\left(\delta_{1}+\delta_{N}+\delta_{2}\right) / t-1\right\}\right] \tag{50}
\end{align*}
$$

## 7. Cases with electrode overlap at the outside of the coil pack

Referring to Fig. 10 it can be seen that there is void space between the coil and the can in which further positive or negative electrode could be incorporated.


Fig. 12. Cases 1 and 4. Filling of void between pack and can (a) by positive and (b) by negative. Cases 2 and 3. Filling of void between pack and can (c) by positive and (d) by negative. Key as for Fig. 10.

### 7.1. Cases 1 and 4

The two possibilities for filling the outer void space are shown in Figs. 12(a) and 12(b).

In the analysis presented earlier (eqns. 15-25) $r$ may be expressed in terms of $r_{c}, \delta_{p}, \delta_{1}, \delta_{N}$ and $\delta_{2}$ for the negative electrode:

$$
\begin{equation*}
r=r_{c}+(1-a)\left(\delta_{p}+\delta_{1}\right)+\delta_{N} / 2 \tag{51}
\end{equation*}
$$

## Positive filling void space

The gap between the pack and the can which may be filled by separators 1 and 2 plus the positive electrode is equivalent to EJ in Fig. 2. The maximum penetration of the positive into the void space ( $\lambda_{m}-\alpha$ ) results when:

$$
\begin{equation*}
\mathbf{E J}=\delta_{1}+\delta_{p}+\delta_{2} \quad\left(\lambda=\lambda_{m}\right) \tag{52}
\end{equation*}
$$

$\lambda_{m}-\alpha$ may be determined by solution of eqns. (52) and (26). The maximum length of overlap by the positive ( $L_{p}^{\prime \prime}$ ) can be calculated using eqn. (2):

$$
\begin{equation*}
L_{p}^{\prime \prime}=\left(\lambda_{m}-\alpha\right)\left\{r_{c}+\delta_{p} / 2-a\left(\delta_{1}+\delta_{p}\right)+n t+\left(\lambda_{m}-\alpha\right) t / 4 \pi\right\} \tag{53}
\end{equation*}
$$

Similarly the corresponding length of separator 1 is $L_{1}^{\prime \prime}$ where:

$$
\begin{equation*}
L_{1}^{\prime \prime}=\left(\lambda_{m}-\alpha\right)\left\{r_{c}+\delta_{p}+\delta_{1} / 2-a\left(\delta_{1}+\delta_{p}\right)+n t+\left(\lambda_{m}-\alpha\right) t / 4 \pi\right\} \tag{54}
\end{equation*}
$$

And the corresponding length of separator 2 is $L_{2}^{\prime \prime}$ where:

$$
\begin{equation*}
L_{2}^{\prime \prime}=\left(\lambda_{m}-\alpha\right)\left\{r_{c}-\delta_{2} / 2-a\left(\delta_{p}+\delta_{1}\right)+n t+\left(\lambda_{m}-\alpha\right) t / 4 \pi\right\} \tag{55}
\end{equation*}
$$

Negative filling void space
The relevant gap between the pack and the can which may be filled with negative electrode is EF in Fig. 2. The maximum penetration of the negative into the void space ( $\gamma_{m}$ ) is given by:

$$
\begin{equation*}
\mathrm{EF}=\delta_{N} \quad\left(\gamma=\gamma_{m}\right) \tag{56}
\end{equation*}
$$

$\gamma_{m}$ may be determined by solution of eqns. (56) and (27). The maximum overlap of the negative $L_{N}^{\prime \prime}$ can be calculated as

$$
\begin{equation*}
L_{N}^{\prime \prime}=D_{t} \gamma_{m} / 2 \tag{57}
\end{equation*}
$$

Note that $D_{t}$ here is less than $D_{c}$ by an amount approximately equal to the negative thickness because $D_{t}$ is calculated using eqn. (10) where $r$ is given by eqn. (51), i.e. calculations are based on a plane through the centre of the negative electrode.

### 7.2. Cases 2 and 3

The two possibilities for filling the outer void space are shown in Figs. $12(\mathrm{c})$ and $12(\mathrm{~d})$.

In the earlier analysis (eqns. 15-25) $r$ may be expressed as follows:

$$
\begin{equation*}
r=r_{c}+d\left(\delta_{n}+\delta_{2}\right)+\delta_{p}+\delta_{1} / 2 \tag{58}
\end{equation*}
$$

## Positive filling void space

The gap between the pack and the can which may be filled by separator 1 plus the positive electrode is equivalent to EF in Fig. 2. The maximum penetration of the positive into the void space $\left(\gamma_{m}\right)$ is given by:

$$
\begin{equation*}
\mathrm{EF}=\delta_{1}+\delta_{p} \quad\left(\gamma=\gamma_{m}\right) \tag{59}
\end{equation*}
$$

$\gamma_{m}$ may be determined by solution of eqns. (59) and (27). The maximum overlap of the positive $L_{p}^{\prime \prime}$ can be calculated as:

$$
\begin{equation*}
L_{p}^{\prime \prime}=\left(D_{t}-\delta_{1}-\delta_{p}\right) \gamma_{m} / 2 \tag{60}
\end{equation*}
$$

Similarly the corresponding length of separator 1 is $L_{1}^{\prime \prime}$ where:

$$
\begin{equation*}
L_{1}^{\prime \prime}=D_{t} \gamma_{m} / 2 \tag{61}
\end{equation*}
$$

## Negative filling void space

The relevant gap between the pack and the can which may be filled by the negative electrode is EJ in Fig. 2. The maximum penetration into the void space ( $\lambda_{m}-\alpha$ ) is given by:

$$
\begin{equation*}
\mathrm{EJ}=\delta_{N} \quad\left(\lambda=\lambda_{m}\right) \tag{62}
\end{equation*}
$$

$\lambda_{m}$ may be determined by solution of eqns. (62) and (26). The maximum overlap of the negative $L_{N}^{\prime \prime}$ can be calculated using eqn. (2) as:

$$
\begin{equation*}
L_{N}=\left(\lambda_{m}-\alpha\right)\left\{r_{c}+\delta_{N} / 2-h\left(\delta_{2}+\delta_{N}\right)+n t+\left(\lambda_{m}-\alpha\right) t / 4 \pi\right\} \tag{63}
\end{equation*}
$$

## 8. Calculation of total component lengths

The steps used in the calculations were as follows:
(i) Determine $n$ using eqn. (34) for one of cases 1-4.
(ii) Calculate $L_{p}, L_{1}, L_{N}$ and $L_{2}$ using eqns. (29) - (32).
(iii) As steps (ii) - (vii) of Section 3.1 using eqns. (51) and (58) for $r$.
(iv) Calculate the partial component lengths $L_{p}^{\prime},\left(L_{1}^{\prime}\right)_{p},\left(L_{2}^{\prime}\right)_{p}, L_{N}^{\prime},\left(L_{1}\right)_{N}$ and $\left(L_{2}^{\prime}\right)_{N}$ using the equations in Section 6 for overlap at the centre of the coil pack.
(v) Calculate the partial component lengths $L_{p}^{\prime \prime}, L_{1}^{\prime \prime}, L_{2}^{\prime \prime}$ and $L_{N}^{\prime \prime}$ using the equations in Section 7 for overlap at the outside of the coil pack.
(vi) The maximum total component lengths may be determined by adding up the calculated partial lengths as follows:
(a) Positive filling void spaces at centre and outside of pack

$$
\begin{aligned}
& \bar{L}_{p}=L_{p}+L_{p}^{\prime}+L_{p}^{\prime \prime} \\
& \bar{L}_{N}=L_{N} \\
& \bar{L}_{1}=L_{1}+\left(L_{1}^{\prime}\right)_{p}+L_{1}^{\prime \prime} \\
& \bar{L}_{2}=L_{2}+\left(L_{2}^{\prime}\right)_{p}+L_{2}^{\prime \prime}
\end{aligned}
$$

(b) Positive filling centre void space, negative filling void between pack and can

$$
\begin{aligned}
& \bar{L}_{p}=L_{p}+L_{p}^{\prime} \\
& \bar{L}_{N}=L_{N}+L_{N}^{\prime \prime} \\
& \bar{L}_{1}=L_{1}+\left(L_{1}^{\prime}\right)_{p} \\
& \bar{L}_{2}=L_{2}+\left(L_{2}^{\prime}\right)_{p}
\end{aligned}
$$

(c) Negative filling centre void space: positive filling void between pack and can

$$
\bar{L}_{p}=L_{p}+L_{p}^{\prime \prime}
$$

$$
\begin{aligned}
& \bar{L}_{N}=L_{N}+L_{N}^{\prime} \\
& \bar{L}_{1}=L_{1}+\left(L_{1}^{\prime}\right)_{N}+L_{1}^{\prime \prime} \\
& \bar{L}_{2}=L_{2}+\left(L_{2}^{\prime}\right)_{N}+L_{2}
\end{aligned}
$$

(d) Negative filling void spaces at the centre and outside of the pack

$$
\begin{aligned}
& \bar{L}_{p}=L_{p} \\
& \bar{L}_{N}=L_{N}+L_{N}^{\prime}+L_{N}^{\prime \prime} \\
& \bar{L}_{1}=L_{1}+\left(L_{1}^{\prime}\right)_{N} \\
& \bar{L}_{2}=L_{2}+\left(L_{2}^{\prime}\right)_{N}
\end{aligned}
$$

## Conclusion

Equations are presented from which the component lengths may be calculated for the four cases corresponding to the permutations of the two electrodes starting at the inside and ending at the outside. These result in gaps between the mandrel and the coil pack and the latter and the can which may be filled with active material for maximized cell capacity. This leads to four variations on the basic cases, resulting in 16 combinations altogether.

The calculations may be readily performed on a programmable calculator or computer.

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## Reference

1 N. Marincic, J. Appl. Electrochem., 6 (1976) 51.

Nomenclature
$A, B$
$a, d, h, k$
$D_{c}$
$D_{t}$
EF, EJ
K
L
constants in eqn. (27),
see Table 4, specified internal can diameter containing coil pack, minimum circle diameter containing spiral, see Fig. 2,
$=2 \pi(n-1 / 2)+t / 2 \pi(r+n t)$,
length of spiral,

| $L_{1}, L_{2}, L_{p}, L_{N}$ | length of separator 1 , separator 2 , positive, negative, with no overlap at centre or between pack and can, |
| :---: | :---: |
| $L_{p}^{\prime},\left(L_{1}^{\prime}\right)_{p},\left(L_{2}^{\prime}\right)_{p}$ | length of positive protruding into the centre void space and the associated lengths of separators 1 and 2, |
| ${ }_{N}^{\prime},\left(L_{1}^{\prime}\right)_{N},\left(L_{2}^{\prime}\right)_{N}$ | as above but for negative, |
| $L_{p}^{\prime \prime}, L_{1}^{\prime \prime}, L_{2}^{\prime \prime}$ | length of positive protruding into the void space between pack and can and the associated lengths of separators 1 and 2, |
| $L_{N}^{\prime \prime}$ | length of negative protruding into the void space between pack and can, |
| $\bar{L}_{1}, \bar{L}_{2}, \bar{L}_{p}, \bar{L}_{N}$ | maximum total length of separator 1 , separator 2 , positive, negative, |
| $n$ | total number of turns on spiral; number of turns of the component at the outside of the coil pack, |
| $n^{\prime}$ | number of turns to C (Fig. 2), |
| $r$ | radius of circle on which spiral originates, |
| $r_{c}$ | radius of mandrel on which coil is wound, |
| $t$ | distance by which spiral moves out during one turn, |
| $\alpha, \gamma, \theta, \lambda$ | see Fig. 3, |
| $\beta$ | angle $\angle \mathrm{ACO}$ (Fig. 2), |
| $\gamma_{m},\left(\lambda_{m}-\alpha\right)$ | angles corresponding to the maximum penetration of an electrode into void space between pack and can, |
| $\epsilon$ | angle $\angle O A C$ (Fig. 2), |
| $\delta_{p}, \delta_{N}$ | thickness of positive, negative, |
| $\delta_{1}, \delta_{2}$ | compressed thickness of separator 1,2 , |
| $\rho_{n}$ | distance of spiral from centre after $n$ turns. |

## Appendix

Derivation of eqn. (4)
The relevant part of Fig. 2 is shown in Fig. 13. A line has been drawn making an angle $\Delta \eta$ with $O C$ intersecting the spiral at $R$.

$$
\begin{align*}
& \angle \mathrm{TCS}=\angle \mathrm{OCP}=\beta  \tag{64}\\
& \mathrm{RO}=r+\left(n^{\prime}+\frac{\Delta \eta}{2 \pi}\right) t  \tag{65}\\
& \mathrm{OS}=\left(r+n^{\prime} t\right) \sec \Delta \eta  \tag{66}\\
& \mathrm{RS}=\mathrm{RO}-\mathrm{OS}=r+\left(n^{\prime}+\frac{\Delta \eta}{2 \pi}\right) t-\left(r+n^{\prime} t\right) \sec \Delta \eta  \tag{67}\\
& \mathrm{SC} \quad=\left(r+n^{\prime} t\right) \tan \Delta \delta \tag{68}
\end{align*}
$$

Consider now the situation $\Delta \eta \rightarrow 0$ only. Then eqn. (67) becomes:


Fig. 13. A section of the spiral required for the derivation of eqn. (4).
Fig. 14. The segment RSC for $\Delta \eta \rightarrow 0$.

$$
\begin{equation*}
\mathrm{RS}=\Delta \eta t / 2 \pi \tag{69}
\end{equation*}
$$

as sec $\Delta \eta=1+(\Delta \eta)^{2} / 2+$ higher order terms and eqn. (68) becomes:

$$
\begin{equation*}
\mathrm{SC}=\left(r+n^{\prime} t\right) \Delta \eta \tag{70}
\end{equation*}
$$

as $\tan \Delta \eta=\Delta \eta+$ higher order terms. The segment RSC in Fig. 13 may then be drawn as shown in Fig. 14. Clearly

$$
\begin{equation*}
\tan \beta=t / 2 \pi\left(r+n^{\prime} t\right) \tag{4}
\end{equation*}
$$

